

Generalised splines via the maximum principle

Andrew D. Lewis*

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*DEPARTMENT OF MATHEMATICS AND STATISTICS, QUEEN'S UNIVERSITY
EMAIL: ANDREW.LEWIS@QUEENSU.CA
URL: [HTTP://WWW.MAST.QUEENSU.CA/~ANDREW/](http://www.mast.queensu.ca/~andrew/)

1. The approach

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- *Question:* What is a spline?
- *My answer:* A curve satisfying a differential equation arising from a minimisation problem.
- Typically, the necessary conditions arising from the minimisation problem are derived with a variational approach.
- Instead, I will use the maximum principle.
- This allows the solution of more general minimisation problems, including, for example, control constraints.
- The control systems I employ are well-suited to the generation of wide classes of curves on manifolds: *affine connection control systems*.

2. What are affine connection control systems?

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- Shortly, they are this:
 1. a configuration manifold Q ;
 2. an affine connection ∇ on Q ;
 3. a collection $\mathcal{Y} = \{Y_1, \dots, Y_m\}$ of vector fields on Q .
- The corresponding control system is

$$\nabla_{c'(t)} c'(t) = u^a(t) Y_a(c(t))$$

for a controlled trajectory (u, c) .

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- Mechanical examples of affine connection control systems:
 1. Lagrangian systems with kinetic energy Lagrangians (∇ is the Levi-Civita connection for the kinetic energy Riemannian metric). For example (some of these need potential energy),
 - satellites,
 - robotic manipulators,
 - underwater vehicles, etc.
 2. Same as above with the addition of constraints linear in velocity. For example,
 - locomotion systems (wheeled vehicles),
 - grasping applications, etc.

3. Affine connection control systems as control affine systems

- Convert

$$\nabla_{c'(t)} c'(t) = u^a(t) Y_a(c(t))$$

to control affine system on TQ :

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$$\dot{v}(t) = f_0(v(t)) + u^a(t) f_a(v(t)),$$

$v \in TQ$.

- Turns out that
 1. the drift is the geodesic spray denoted $f_0 = Z$, and
 2. the control vector fields are the vertical lifts of the vectors fields from \mathcal{Y} : we write $f_a = Y_a^{\text{lift}}$.

4. The Maximum Principle for affine connection control systems

- Noakes, Heinzinger, Paden, and Camarinha, Crouch, Silva Leite, and Sontag, Sussmann, and Fax, Murray, and Chyba, Leonard, Sontag.
- We shall investigate in a little detail *one* of the several consequences of the Maximum Principle as it applies to affine connection control systems.

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- Start general—let's look at the Maximum Principle for

$$c'(t) = f_0(c(t)) + u^a(t) f_a(c(t)),$$

with $c(t) \in M$, u taking values in $U \subset \mathbb{R}^m$, and objective function $L(x, u)$.

- Have the **control Hamiltonian** on $U \times T^*M$:

$$H(\alpha_x, u) = \underbrace{\alpha_x(f_0(x))}_{H_1} + \underbrace{\alpha_x(u^a f_a(x))}_{H_2} - \underbrace{L(x, u)}_{H_3}.$$

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- One of several consequences of the MP is that if (u, c) is a minimiser then there exists a one-form field λ along c with the property that $t \mapsto \lambda(t)$ is an integral curve for the time-dependent Hamiltonian $(\alpha_x, t) \mapsto H(\alpha_x, u(t))$.
- The Hamiltonian is a sum of three terms, and so too will be the Hamiltonian vector field. Let us look at the first term, that with (plain old) Hamiltonian $H_1(\alpha_x) = \alpha_x(f_0(x))$.
- In local coordinates X_{H_1} is written as

$$\begin{aligned} \dot{x}^i &= f_0^i(x) \\ \dot{p}_i &= -\frac{\partial f_0^j}{\partial x^i} p_j \quad \longleftarrow \quad \text{"adjoint equation"?'} \end{aligned}$$

- X_{H_1} is the **cotangent lift** of f_0 and we denote it $f_0^{T^*}$.

- **Objective:** Understand $f_0^{T^*}$ when $M = TQ$ and $f_0 = Z$.
- Begin with a change of subject: Let f_0 be a vector field on (general) M with f_0^T its **tangent lift** defined by

$$f_0^T(v_x) = \left. \frac{d}{dt} \right|_{t=0} T_x F_t(v_x)$$

(F_t is the flow of f_0).

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- f_0^T is the "linearisation" of f_0 and in coordinates is given by

$$\begin{aligned} \dot{x}^i &= f_0^i(x) \\ \dot{v}^i &= \frac{\partial f_0^i}{\partial x^j} v^j \end{aligned} \quad \left(\begin{array}{l} \text{compare } f_0^{T^*} : \\ \dot{x}^i = f_0^i(x) \\ \dot{p}^i = -\frac{\partial f_0^j}{\partial x^i} p_j \end{array} \right)$$

- The flow of f_0^T measures how the integral curves of f_0 change as we change the initial condition in the direction of v_x .

- Perhaps we can understand Z^T —thus take $M = TQ$ and $f_0 = Z$ in the discussion of tangent lift.
- Note:
 - Projections of integral curves of Z to Q are geodesics of ∇ .
 - Z^T measures variations of integral curves of Z .
 - Thus Z^T measures variations of geodesics.

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- But we know something else which measures variations of geodesics. . .
- Let $c(t)$ be a geodesic. By varying the initial condition for the geodesic we generate an “infinitesimal variation” ξ of the geodesic and it turns out to satisfy. . . the **Jacobi equation**:

$$\nabla_{c'(t)}^2 \xi(t) + R(\xi(t), c'(t))c'(t) + \nabla_{c'(t)}(T(\xi(t), c'(t))) = 0.$$

- What is the *precise* relationship between Z^T and the Jacobi equation?

Some tangent bundle geometry using Z

- To make the “connection” between Z^T and the Jacobi equation, we perform constructions on the tangent bundle using the spray Z .
- ∇ comes from a linear connection on Q which induces an Ehresmann connection on $\pi_{TQ}: TQ \rightarrow Q$.
- Thus we may write $T_{v_q}TQ \simeq T_qQ \oplus T_qQ$.
- Z^T is not a spray, but. . . if $I_Q: TTQ \rightarrow TTQ$ is the canonical involution then $I_Q^*Z^T$ is a spray (it is the spray for the so-called *complete lift of ∇*).
- Use $I_Q^*Z^T$ to induce an Ehresmann connection on $\pi_{TTQ}: TTQ \rightarrow TQ$.
- Thus

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$$\begin{aligned} T_{X_{v_q}}TTQ &\simeq T_{v_q}TQ \oplus T_{v_q}TQ \\ &\simeq \underbrace{T_qQ \oplus T_qQ}_{\text{geodesic equations}} \oplus \underbrace{T_qQ \oplus T_qQ}_{\text{variation equations}} \end{aligned}$$

- One represents Z^T in this splitting and determines that the Jacobi equation sits “inside” one of the four components.
- Now one applies similar constructions to T^*TQ and Z^{T^*} to derive (all going to plan) a one-form version of the Jacobi equation.
- Need a little notation:

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$$\langle R^*(\alpha, u)v; w \rangle = \langle \alpha; R(w, u)v \rangle, \quad \langle T^*(\alpha, u); w \rangle = \langle \alpha; T(w, u) \rangle.$$

- After the dust settles, we get what we are after which is the **adjoint Jacobi equation**:

$$\nabla_{c'(t)}^2 \lambda(t) + R^*(\lambda(t), c'(t))c'(t) - T^*(\nabla_{c'(t)} \lambda(t), c'(t)) = 0.$$

- Why did I do this?
 - The adjoint Jacobi equation captures the interesting part of the Hamiltonian vector field Z^{T^*} , which comes from the MP, and words it in terms of affine differential geometry, i.e.,



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$$\begin{aligned} \nabla_{c'(t)} c'(t) &= 0 \\ \nabla_{c'(t)}^2 \lambda(t) + R^*(\lambda(t), c'(t))c'(t) - T^*(\nabla_{c'(t)} \lambda(t), c'(t)) &= 0. \end{aligned}$$

- The geometry of Z on TQ provides a way of **globally** pulling out the “adjoint equation” from the MP in an intrinsic manner—this is not generally possible in the MP.

- The adjoint Jacobi equation forms the backbone of a general statement of the MP for affine connection control systems.
 - The contribution of the inputs needs to be added (easy).
 - The contribution of the objective function needs to be added (difficulty depends on the nature of the function).

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- Take the case when objective function is $L(u, v_q) = \frac{1}{2}g(v_q, v_q)$ for a Riemannian metric, and the affine connection is not necessarily the Levi-Civita connection. (In the case when ∇ is the Levi-Civita connection, a result is obtained by Silva Leite, Camarinha, and Crouch.)
- In this case, it is possible for there to be abnormal extremals (and probably abnormal minimisers).

- The normal extremals satisfy

$$\begin{aligned} \nabla_{c'(t)} c'(t) &= -h_Y^\sharp(\lambda(t)) \\ \nabla_{c'(t)}^2 \lambda(t) + R^*(\lambda(t), c'(t))c'(t) - T^*(\nabla_{c'(t)} \lambda(t), c'(t)) &= \\ &= \frac{1}{2} \nabla h_Y(\lambda(t), \lambda(t)) - T^*(\lambda(t), h_Y^\sharp(\lambda(t))), \end{aligned}$$

and abnormal extremals satisfy three conditions:

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1. $\nabla_{c'(t)} c'(t) = u^a(t) Y_a(c(t))$,
2. $\lambda(t) \in \text{ann}(Y_{c(t)})$ for $t \in [a, b]$ and
3. λ satisfies the equation along c given by:

$$\begin{aligned} \nabla_{c'(t)}^2 \lambda(t) + R^*(\lambda(t), c'(t))c'(t) - T^*(\nabla_{c'(t)} \lambda(t), c'(t)) &= \\ &= B_Y(\lambda(t), u^a(t) Y_a(t)). \end{aligned}$$

- When ∇ is the Levi-Civita connection for g , and when the system is fully actuated, then we recover the equation of Noakes, Heinzinger, and Paden and Crouch and Silva Leite:

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$$\nabla_{c'(t)}^3 c'(t) + R(\nabla_{c'(t)} c'(t), c'(t)) = 0.$$

- Where to go from here?
 - Other cost functions (time, length, etc.)
 - Constructible examples.